

# Plenoptic Principal Planes

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**Abstract:** We show that the plenoptic camera is optically equivalent to an array of cameras. We compute the parameters that establish that equivalence and show where the plenoptic camera is more useful than the camera array.

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## 1. Introduction

In 1908, Lippmann introduced the idea of using a microlens array to capture the radiance of a scene and produce what he called *integral photographs* [1]. These photographs displayed not just a 2D picture, but also “the 3D relief” of the scene [2]. In 1928, Ives subsequently added an objective lens to Lippmann’s microlens array [3] to increase the acuity. Since then, many other researchers and scientists have continued to advance integral photography.

With the advent of digital photography, significant new opportunities to investigate integral photography became available. In 1992, Adelson introduced the “plenoptic camera,” a digital version of Lippmann’s microlens array designed to solve problems in computer vision [4]. In 2005, Ng improved the plenoptic camera and introduced new digital-processing methods, including refocusing [5, 6]. In later work Fife [7] and Lumsdaine and Georgiev [8] improved the resolution by introducing the focused plenoptic camera and developed methods for rendering synthetic views with computational focus at high resolution.

In parallel with this, researchers have also used arrays of ordinary cameras to capture multiple views, reconstruct the plenoptic function, and render synthetic refocused views [9]. Designs using a single camera, but with arrays of lenses in front of the main lens have also been proposed [10].

Thus, we have two families of systems for capturing the plenoptic function: those with a single main lens and a microlens array and those with multiple cameras (multiple main lenses). A natural question then arises regarding these two families. Are these two families equivalent in some sense or are they fundamentally different? In this paper we prove that these two families are in fact optically equivalent and we derive the parameters that establish this correspondence.

## 2. Principal Planes in Integral Photography

In ray optics light transport is described by the ray matrices of the corresponding optical elements [11]. These matrices act on the ray vector  $(q, p)$ , where  $q$  is the displacement away from the optical axis,  $p$  is the slope. Since in Gaussian optics the system is linear, the joint action of all elements is described by the product of their matrices in the order in which light encounters them when travelling through an optical system.

A lens of focal length  $f$  and a translation of distance  $t$  are respectively described by the ray matrices

$$L(f) = \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix}, \quad T(t) = \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix}. \quad (1)$$

As we see in Figure 1, the three optical elements that light encounters while travelling in a plenoptic camera are the main lens (of focal length  $F$ ), the space  $D$  from the main lens to the microlens array, and the typical microlens (of focal length  $f$ ). The ray matrix of this system would be

$$L(f)T(D)L(F) = \begin{bmatrix} 1 - \frac{D}{F} & D \\ \frac{D}{fF} - \frac{1}{f} - \frac{1}{F} & 1 - \frac{D}{f} \end{bmatrix}. \quad (2)$$

In order to derive expressions for the location and focal lengths of the virtual lenses, the well-known concept of principal planes (due to Gauss) will be used [12]. With principal planes, any optical system made of lenses and spaces

between them can be replaced with an optically equivalent system consisting of space from the entrance to the “first principal plane”; followed by refraction by appropriate lens; followed by travel from a “second principal plane” to the exit. Mathematically, this means that the matrix corresponding to any optical system can be formed as the product of three matrices:

$$T(y)L(\Phi)T(x) = \begin{bmatrix} 1 & y \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -\frac{1}{\Phi} & 1 \end{bmatrix} \begin{bmatrix} 1 & x \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 - \frac{y}{\Phi} & x + y - \frac{xy}{\Phi} \\ -\frac{1}{\Phi} & 1 - \frac{x}{\Phi} \end{bmatrix}. \quad (3)$$

Note that the distances  $x$  and  $y$  don't have to add up to the total distance travelled, and the principal planes may actually be outside the system (in this case  $x$  and/or  $y$  are negative).

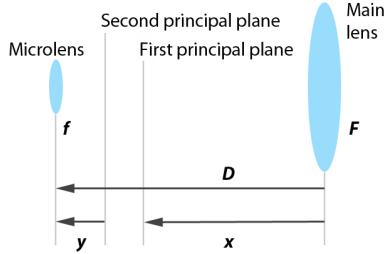


Fig. 1: The concept of principal planes applied to the plenoptic camera. Only one microlens is shown.

Now we equate the system matrix described by (2) with the principal planes system matrix (3). By making the bottom left matrix element in (2) and (3) equal, we derive the focal length  $\Phi$  of the effective lens. Similarly we derive the distances  $x$  and  $y$ .

$$\Phi = \frac{Ff}{F + f - D}, \quad x = \frac{DF}{f + F - D}, \quad y = \frac{Df}{f + F - D}. \quad (4)$$

For the incoming light the virtual lens is positioned at the first principal plane; for the sensor the virtual lens is positioned at the second principal plane. It is well known that the focused plenoptic camera exists in two versions: As an array of Keplerian or Galilean telescopes. Distances and focal lengths are positive or negative depending on which type of focused plenoptic camera we are using, with real or virtual image of the main lens image.

### 3. Cameras

A typical focused plenoptic camera is described in Lumsdaine and Georgiev [8]. We will establish the exact equivalence of this camera with an appropriate camera array. The main lens has focal length  $F = 140mm$ . All microlenses have focal length  $f = 750\mu m$ . We consider the two types of camera separately.

#### 3.1. Keplerian case

The value of  $D$  can be freely modified by the user in such a way that the denominator  $F + f - D$  actually goes through zero. To get a reasonable value of  $D$ , let's assume that the world is at infinity, and that the main lens image is formed at a distance of  $6f$  in front of the microlens. Such settings are typical for a real camera. Then a simple calculation shows that the focal length  $\Phi = -28mm$ , and the distances are  $x = -5226mm$ , and  $y = -28mm$ , approximately. Figure 2 shows the dependence of the focal length of the virtual lens on the value of  $D$ . The two branches of the graph correspond to the Keplerian and Galilean cases. The specific value of  $\Phi = -28mm$  is represented as a dot.

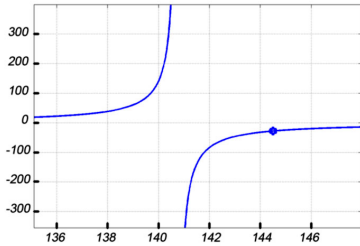


Fig. 2: Focal length of the virtual lens as a function of the distance  $D$  between main lens and microlens.

We see that the equivalent virtual lens is negative, i.e. the image is not flipped. That's what we actually observe in the captured imagery: Microimages are not flipped relative to the world. Flipping two times in the relay system is equal to no flip. More interesting, the virtual lens is placed at big negative  $x$ , far in front of the camera. This is as if the plenoptic camera has an array of negative lenses in front, very similar to the camera described in [10]. The distance  $y$  is also negative. That's natural considering that the image of a negative lens is actually formed in front of it. It is also as expected that the image of an object at infinity is formed at one focal length from the lens.

The distance between microlenses is  $d = 250\mu m$ . Considering similar triangles in the camera geometry, the distance between the virtual lenses that are in front of the main lens would be  $xd/D = 9.3mm$ . Again, we observe approximately same size as that of the negative lenses that were placed in front of the camera in the case of [10].

### 3.2. Galilean case

To fix the value of  $D$  let's assume that the world is at infinity, and the main lens image is formed at a distance of  $6f$  behind the microlens. Then a simple calculation shows that the focal length  $\Phi = 28mm$ , and the distances are  $x = 5226mm$ , and  $y = 28mm$ , approximately. We see that the virtual lens is positive and the image would be flipped. That's what we actually observe in the captured imagery. The virtual lens is placed at big positive  $x$ , far behind the camera (as in a Galilean telescope). The distance  $y$  is also positive. It is as expected that the image of an object at infinity is formed at one focal length behind the lens.

In both the Keplerian and Galilean cases the difference from a camera array is that now the lenses are virtual, and we can actually position them within a big range, including locations that are mechanically not accessible. Also, our equations are very sensitive to the value of the denominator  $F + f - D$  so we can easily tune and vary the distances within a wide range, including infinity.

## 4. Conclusion

In this paper we have presented a proof that the two approaches to radiance capture, the plenoptic and the multicamera array, are equivalent. We establish the correspondence between the plenoptic camera and the array of cameras based on the concept of principal planes, and we compute the parameters for that correspondence in two typical cases.

In our opinion the two main factors that make the plenoptic camera preferable to the camera array are the smaller size, and the lower cost of a quality plenoptic implementation – compared to the corresponding equivalent camera array. Also, the great flexibility in the positioning the virtual lenses in the plenoptic camera can be practically useful.

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